



## TROSY NMR with partially deuterated proteins

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### Abstract

TROSY-type optimization of liquid-state NMR experiments is based on the preservation of unique coherence transfer pathways with distinct transverse relaxation properties. The broadband decoupling of the <sup>1</sup>H spins interchanges the TROSY and anti-TROSY magnetization transfer pathways and thus is not used in TROSY-type triple resonance experiments or is replaced with narrowband selective decoupling. To achieve the full advantage of TROSY, the uniform deuteration of proteins is usually required. Here we propose a new and general method for <sup>1</sup>H broadband decoupling in TROSY NMR, which does not compromise the relaxation optimization in the <sup>15</sup>N-<sup>1</sup>H moieties, but uniformly and efficiently refocuses the <sup>1</sup>J<sub>CH</sub> scalar coupling evolution in the <sup>13</sup>C-<sup>1</sup>H moieties. Combined with the conventional <sup>2</sup>H decoupling, this method enables obtaining high sensitivity TROSY-type triple resonance spectra with partially deuterated or fully protonated <sup>13</sup>C, <sup>15</sup>N labeled proteins.

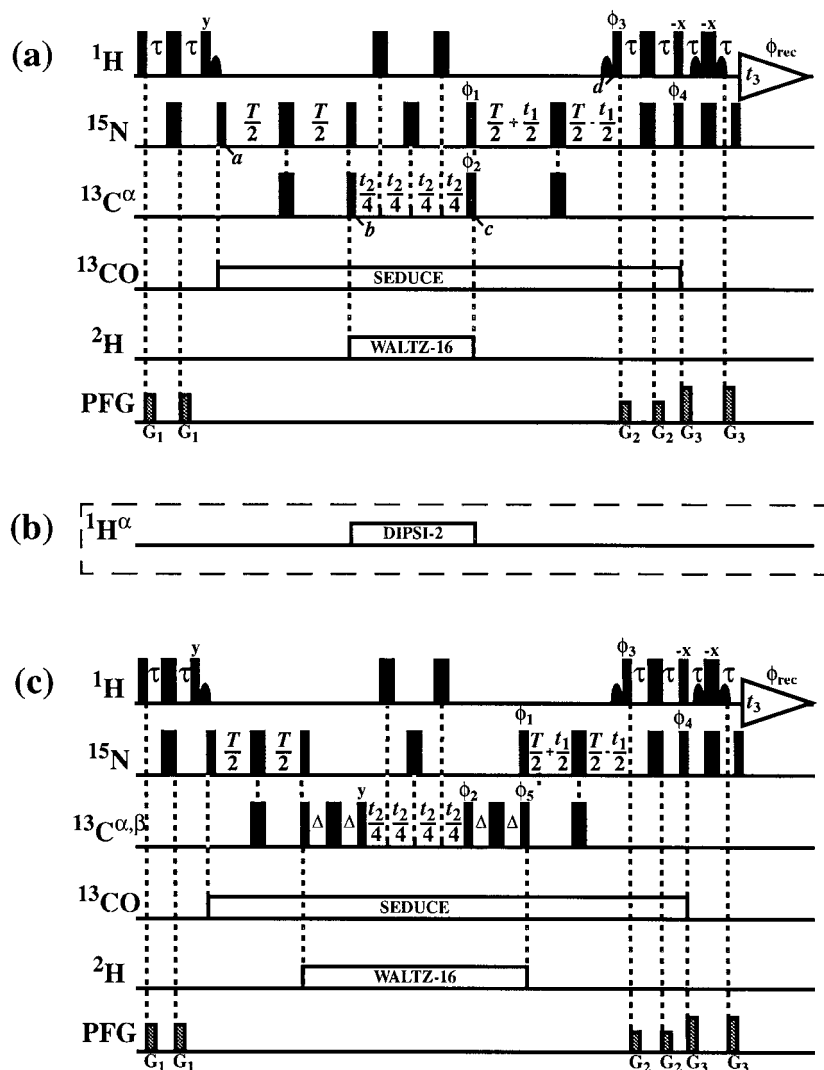
*Abbreviations:* TROSY, transverse relaxation-optimized spectroscopy; DD, dipole–dipole coupling; CSA, chemical shift anisotropy; 2D, two-dimensional.

Uniform or partial replacement of nonlabile protons with deuterons significantly reduces transverse relaxation by scaling down dipole–dipole interactions with remote hydrogens (LeMaster, 1994; Yamazaki et al., 1994; Grzesiek et al., 1995; Shan et al., 1996; Gardner and Kay, 1998). In all cases transverse relaxation of the amide protons and remaining aliphatic and aromatic protons is substantially reduced (LeMaster, 1994; Gschwind et al., 1999). Transverse relaxation-optimized spectroscopy (Pervushin et al., 1997) is shown to be effective in further improving sensitivity of the triple resonance experiments performed with both protonated and uniformly deuterated <sup>13</sup>C, <sup>15</sup>N labeled proteins (Salzmann et al., 1998). TROSY-type optimization of NMR experiments is based on the preservation of the unique magnetization transfer pathways with distinct transverse relaxation properties (Pervushin et al., 1997). The broadband decoupling of the <sup>1</sup>H spins interchanges the TROSY and

anti-TROSY coherence transfer pathways and thus is avoided in the TROSY-type triple resonance experiments or is replaced with a narrowband selective decoupling (Salzmann et al., 1998). To achieve the full advantage of TROSY the uniform deuteration of proteins was usually required. We propose a new and general method for the <sup>1</sup>H broadband decoupling in TROSY NMR, which does not compromise the relaxation optimization in the <sup>15</sup>N-<sup>1</sup>H moieties, but uniformly and efficiently refocuses the <sup>1</sup>J<sub>CH</sub> scalar coupling evolution in the <sup>13</sup>C-<sup>1</sup>H moieties. This method makes the TROSY-type triple resonance experiments suitable for obtaining backbone resonance assignment of partially deuterated and fully protonated <sup>13</sup>C, <sup>15</sup>N labeled proteins.

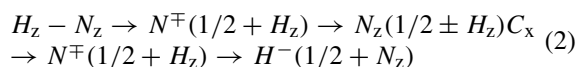
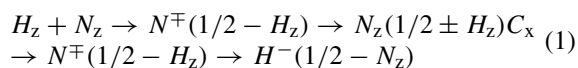
The two relevant magnetization transfer pathways in TROSY-type triple resonance experiments HNCA and HNCACB shown in Figure 1a and c with the most and the least favorable transverse relaxation properties can be schematically represented as (Pervushin et al.,

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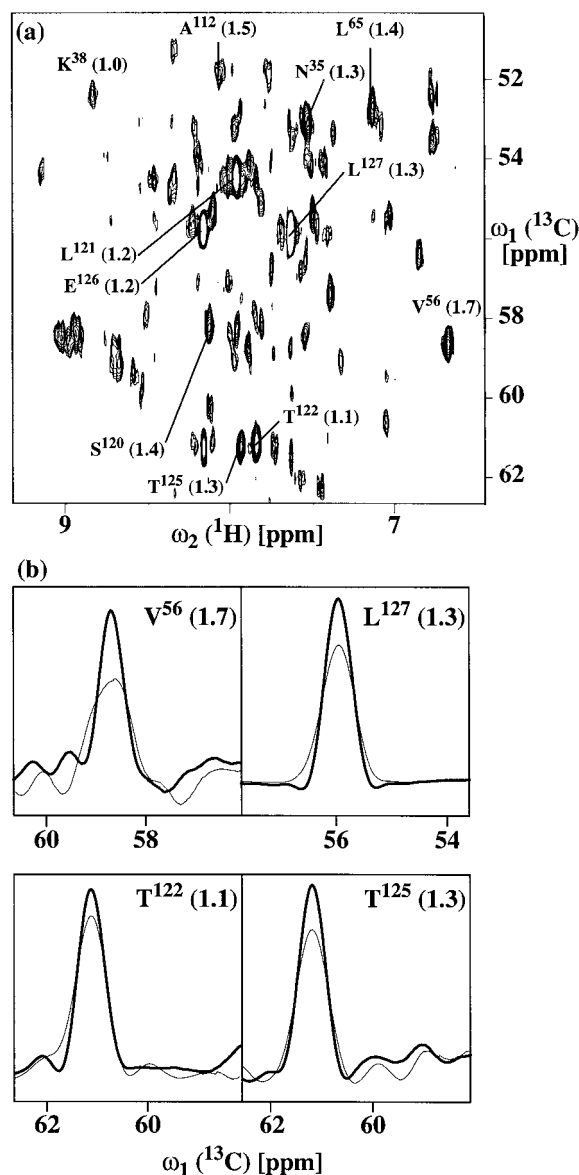
**Figure 1.** Experimental scheme for the  $[^{15}\text{N}, ^1\text{H}]$ -TROSY-HNCA (a) and  $[^{15}\text{N}, ^1\text{H}]$ -TROSY-HNCACB (c) experiments. In (a) the radio-frequency pulses on  $^1\text{H}$ ,  $^{15}\text{N}$ ,  $^{13}\text{C}$ ,  $^{13}\text{CO}$ ,  $^2\text{H}$  and  $^1\text{H}^\alpha$  are applied at 4.7, 118, 55, 174, 3.6 and 4.7 ppm, respectively. The same frequencies are used in (c) except for  $^{13}\text{C}$ , which is placed at 46 ppm. Narrow and wide black bars indicate non-selective  $90^\circ$  and  $180^\circ$  pulses, respectively. Sine bell shapes on the line marked  $^1\text{H}$  indicate selective  $90^\circ$  pulses. The line marked PFG indicates the duration and strength of pulsed magnetic field gradients applied along the z-axis:  $G_1$ : 800  $\mu\text{s}$ , 50 G/cm;  $G_2$ : 800  $\mu\text{s}$ , 40 G/cm;  $G_3$ : 800  $\mu\text{s}$ , 70 G/cm. The delays are  $\tau=2.4$  ms and  $T=44$  ms. The phase cycle is:  $\phi_1 = \{y, -y, -x, x\}$ ;  $\phi_2 = \{4x, 4(-x)\}$ ;  $\phi_3 = \{-y\}$ ;  $\phi_4 = \{-y\}$ ;  $\phi_5 = \{4y, 4(-y)\}$ ;  $\phi_{\text{rec}} = \{y, -y, -x, x, -y, y, x, -x\}$ . All other radio-frequency pulses are applied with phase x. A phase-sensitive spectrum in the  $^{15}\text{N}(t_1)$  dimension is obtained by recording a second FID for each  $t_1$  value, with  $\phi_1 = \{y, -y, x, -x\}$ ,  $\phi_3 = \{y\}$  and  $\phi_4 = \{y\}$ , and data processing as described by Kay et al. (1992). Quadrature detection in the  $^{13}\text{C}^\alpha(t_2)$  dimension is achieved by the States-TPPI method (Marion et al., 1989) applied to the phases  $\phi_2$  and  $\phi_5$ . The use of water flip-back pulses (Grzesiek and Bax, 1993) ensures that the water magnetization stays aligned along the +Z axis throughout both the constant-time period  $T$  and the data acquisition period  $^1\text{H}(t_3)$ .  $^2\text{H}$ -decoupling during  $t_2$  is achieved with WALTZ-16 (Shaka et al., 1983) at the field strength of  $\gamma B_2 = 2.5$  kHz. The reference TROSY-HNCA experiment with selective  $^1\text{H}^\alpha$ -decoupling during the  $^{13}\text{C}^\alpha(t_2)$  evolution period (Salzmann et al., 1998) is constructed by performing the DIPSII-2 sequence (Cavanagh and Rance, 1992) with  $\gamma B_2 = 0.51$  kHz as indicated on the insert (b).

1998):



The first arrow designates coherence transfer from the combined  $^1\text{H}$  and  $^{15}\text{N}$  steady state polarizations to the individual transition of the  $^{15}\text{N}$  spin, the second arrow represents the coherence transfer from  $^{15}\text{N}$  to  $^{13}\text{C}$  and the evolution of magnetization on the  $^{13}\text{C}$  spins, and the last two arrows designate the return of magnetization to the individual transition of the  $^1\text{H}$  spin for detection (Salzmann et al., 1998, 1999). The two pathways are distinguished by significantly different relaxation properties of the  $N^\mp(1/2 - H_z)$  and  $N^\mp(1/2 + H_z)$  operators due to the interference between  $^{15}\text{N}$ - $^1\text{H}$  DD and  $^{15}\text{N}$  CSA interactions (Shimizu, 1964; Goldman, 1984) during the coherence transfer phase and by the relaxation properties of the  $H^-(1/2 - N_z)$  and  $H^-(1/2 + N_z)$  operators during signal acquisition. Note that the two  $^{15}\text{N}$  and  $^1\text{H}$  operators and the corresponding magnetization transfer pathways can be mutually interchanged if the longitudinal spin in each operator is inverted by the rf-pulse. Thus, during the coherence transfer periods it is important not to perturb corresponding longitudinal spins, which happens for example during the long constant time periods  $T$  between time points a and b and c and d in the experimental schemes of Figure 1 and during signal acquisition. On the other hand, during the  $^{13}\text{C}$  evolution period between time points b and c, the  $^{15}\text{N}$  magnetization is stored in the form of magnetization modes  $N_z(1/2 + H_z)$  and  $N_z(1/2 - H_z)$ , which can be interchanged without any significant effect on relaxation (Canet, 1989).

In partially deuterated proteins the evolution of the  $^{13}\text{C}$  spins under the scalar coupling to the directly attached hydrogen spin in both  $^{13}\text{C}$ - $^1\text{H}$  and  $^{13}\text{C}$ - $^2\text{H}$  moieties has to be simultaneously refocused. This is achieved by an application of the conventional broadband  $^2\text{H}$  decoupling and two  $180^\circ$   $^1\text{H}$  hard pulses as is shown in Figure 1a and c. Such a combination interchanges the  $^{15}\text{N}$  magnetization modes twice, thus resulting in no mixing of the TROSY and ‘anti-TROSY’ magnetization transfer pathways represented by Equations 1 and 2, respectively. Another advantage of the proposed experimental scheme of Figure 1 is the minimal saturation of the water signal during the  $^1\text{H}^{\alpha,\beta}$  decoupling due to the absence of net rotation of



**Figure 2.** Comparison of broadband  $^1\text{H}$  decoupling versus  $^1\text{H}^\alpha$ -selective narrowband decoupling in TROSY-HNCA experiments. Two spectra were recorded for  $^{15}\text{N}$ ,  $^{13}\text{C}$ -labeled BsCM protein with 30–70% random deuteration free in solution. (a) Contour plot of the broadband  $^1\text{H}$  decoupled 2D [ $^{13}\text{C}$ ,  $^1\text{H}$ ]-correlation projection of the 3D TROSY-HNCA spectrum measured with the experimental scheme of Figure 1a. Values of relative improvement in signal amplitude as compared to the experiment with  $^1\text{H}^\alpha$ -selective decoupling are given in parentheses together with the sequence-specific resonance assignments of the selected intraresidual cross peaks. (b) 1D cross-sections of the 2D TROSY-HNCA spectra along the  $^{13}\text{C}$  dimension at the intraresidual [ $^{13}\text{C}$ ,  $^1\text{H}$ ]-correlation cross peak positions of residues V $^{56}$ , T $^{122}$ , T $^{125}$  and L $^{127}$ . Thick and thin lines correspond to the broadband  $^1\text{H}$ -decoupled and the  $^1\text{H}^\alpha$ -selective decoupled experiments, respectively. The size of the acquired time domain data was  $80 \times 2048$  complex points. Forty-eight scans per increment were accumulated, with a total measuring time of 1 h for both experiments.

the water magnetization between time points b and c. Since the  $^{13}\text{C}$   $t_{2\text{max}}$  delay in both TROSY-HNCA and TROSY-HNCACB experiments is usually limited by  $^1J_{\text{C}\alpha\text{C}\beta}$  coupling to a few milliseconds, the loss of the water signal due to radiation damping between  $180^\circ$   $^1\text{H}$  pulses can be safely neglected (Sobol et al., 1998).

The 44 kDa 30–70% randomly deuterated and uniformly  $^{13}\text{C}$ ,  $^{15}\text{N}$ -labeled water-soluble trimeric enzyme *B. subtilis* Chorismate Mutase (EC 5.4.99.5) (Chook et al., 1994; Ladner et al., 2000) was used in the setup and evaluation of the proposed experiments. Figure 2 compares the 2D [ $^{13}\text{C}$ ,  $^1\text{H}$ ]-correlation projection of the 3D TROSY-HNCA spectrum measured with the experimental scheme of Figure 1a with the corresponding spectrum obtained using the narrowband  $^1\text{H}$  decoupling as described by Salzmann et al. (1998, 1999). A significant improvement of the line shape along the  $^{13}\text{C}$  dimension for the majority of resonances with a concomitant increase of the sensitivity in the range between 1.1 and 1.8 are observed. This is demonstrated in Figure 2b by superimposing 1D  $^{13}\text{C}$  cross-sections taken from the conventional TROSY-HNCA spectrum (thin lines) and the TROSY-HNCA spectrum measured with the experimental scheme of Figure 1a (thick lines) at the positions of the resonances of V<sup>56</sup>, T<sup>122</sup>, T<sup>125</sup> and L<sup>127</sup>.

Overall, the sensitivity improvement of the proposed TROSY-type triple resonance experiments for partially deuterated proteins stems from several factors such as removal of the off-resonance effects of the narrowband  $^1\text{H}^{\alpha,\beta}$  decoupling, which might result in the interchange of the TROSY and ‘anti-TROSY’ pathways, removal of the intrinsic problems of the narrowband decoupling such as long supercycles and arbitrary truncation of the selective pulses due to the time delay incrementation (Matsuo et al., 1996; Van der Kooi et al., 1999), and finally the minimization of the water saturation (Stonehouse et al., 1995). The proposed method of  $^1\text{H}^{\alpha,\beta}$  decoupling further broadens the scope of applicability of the TROSY-type experiments, which renders the TROSY line of experiments an attractive choice for resonance assignment of partially deuterated or fully protonated proteins.

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